

Subgap Kinetic Inductance Detector Sensitive to 85-GHz Radiation

F. Levy-Bertrand^{1,*}, A. Benoît,¹ O. Bourrion², M. Calvo,¹ A. Catalano,² J. Goupy,¹ F. Valenti,^{3,4}
N. Maleeva,^{4,5} L. Grünhaupt⁴, I.M. Pop^{4,6} and A. Monfardini^{1,†}

¹University of Grenoble Alpes, CNRS, Grenoble INP, Institut Néel, 38000 Grenoble, France

²Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble Alpes, CNRS, 53 rue des Martyrs,
38026 Grenoble Cedex, France

³IPE, Karlsruhe Institute of Technology, 76344 Eggenstein-Leopoldshafen, Germany

⁴PHI, Karlsruhe Institute of Technology, 76131 Karlsruhe, Germany

⁵National University of Science and Technology MISIS, 119049 Moscow, Russia

⁶IQMT, Karlsruhe Institute of Technology, 76344 Eggenstein-Leopoldshafen, Germany

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We have fabricated an array of subgap kinetic inductance detectors (SKIDs) made of granular aluminum ($T_c \sim 2$ K) sensitive in the 80–90 GHz frequency band and operating at 300 mK. We measure a noise equivalent power of 1.3×10^{-16} W/Hz^{0.5} on average and 2.6×10^{-17} W/Hz^{0.5} at best, for an illuminating power of 50 fW per pixel. Even though the circuit design of SKIDs is identical to that of the kinetic inductance detectors, the SKIDs operating principle is based on their sensitivity to subgap excitations. This detection scheme is advantageous because it avoids having to lower the operating temperature proportionally to the lowest detectable frequency. The SKIDs presented here are intrinsically selecting the 80–90 GHz frequency band, well below the superconducting spectral gap of the film, at approximately 180 GHz.

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I. INTRODUCTION

Terahertz (THz) electromagnetic waves are defined from 0.1 to 10 THz, corresponding to wavelengths from 3 mm down to 30 μm [1–3]. These waves are nonionizing and can pass through wood, plastic, clothing, and the atmosphere in a few frequency bands (THz atmospheric transmission windows). Most molecules have vibration or rotation frequencies in the THz range, allowing spectroscopic identification [4,5]. These properties drive the development of THz imaging and THz spectroscopy for various potential applications such as security [6], medical diagnostic [7], quality control processes, and astrophysics research [8,9]. To better exploit these potential applications, two main research directions are being conducted: the development of sensitive detectors and the development of intense sources. The present work focuses on the investigation of a sensitive detector, the subgap kinetic inductance detector (SKID) [10], a particular implementation of the kinetic inductance detector (KID) concept [11–13].

KIDs are detectors used for millimeter-wave observations in astrophysics [14–17] and for passive THz cameras emerging for security applications [18,19]. They are planar

resonant circuits made of superconductors deposited on an insulating substrate. They are usually cooled down to about 100 mK. The photon detection principle consists of monitoring the resonance frequency shift that is proportional to the incident power. Thanks to a straightforward fabrication process, an easy multiplexing technique, and a compact design, KIDs are ideal for achieving large sensitive arrays with thousands of pixels, such as those currently installed in the NIKA2 millimetric camera of the IRAM 30 m telescope (Pico Veleta, Spain) [20]. The current NIKA2 KID arrays have a total of about 3000 KIDs fabricated from thin superconducting aluminum, operating at 150 mK, and detecting frequencies in the 120–300 GHz band. Their noise equivalent power is of the order of 10^{-17} to 10^{-16} W/Hz^{0.5} as required for ground-based observations.

The circuit design of SKIDs is identical to that of KIDs. The photon detection principle also consists of monitoring the resonance frequency shift that is proportional to the incident power. However, the mechanism giving rise to the frequency shift is different. For classic KIDs [11], in order to generate a frequency shift the incident photon must carry an energy $h\nu$ higher than or equal to the superconducting spectroscopic gap 2Δ . This gap is the minimum energy required to break the Cooper pairs forming the superfluid. Thus, for a given superconducting material, the smallest detectable frequency is $\nu_{\text{low}} = 2\Delta/h$. An associated

*florence.levy-bertrand@neel.cnrs.fr

†alessandro.monfardini@neel.cnrs.fr

constraint is that, in order to exclude temperature-induced excitations, the operating temperature must be less than or equal to $T \sim T_c/10 \sim 2\Delta/(35k_B)$ where T_c is the superconducting critical temperature and k_B is the Boltzmann constant. In short, for classic KIDs, reducing the smallest detectable frequency requires a proportional reduction of the superconducting gap and, consequently, of the operating temperature. For detection at 80 GHz, an operating temperature lower than ~ 100 mK would thus be needed. The working principle of SKIDs removes this constraint. For SKIDs [10,21], the shift of the measured frequency is due to absorbed photons at a subgap excitation. The frequency shift occurs because, in superconducting thin films, the superfluid density, and thus the kinetic inductance, depends on the circulating current [10,22,23]. The current flowing in the detector is increased by the absorbed photons. For a current I that is small compared with the critical current I_c , the frequency shift is as follows [22,24]:

$$\frac{\delta f}{f} \sim -\frac{\alpha}{2} \left(\frac{I}{I_*} \right)^2 \quad (1)$$

where α is the ratio of the kinetic inductance by the total inductance and $I_* = 3\sqrt{3}/2 \times I_c$. To detect a small incident power, Eq. (1) provides evidence that a small critical current I_c is needed. With this aim, two parameters can be optimized: the critical current density, J_c , that is material dependent, and the section of the kinetic inductance design. This effect can also be formalized in the language of circuit quantum electrodynamics, using the cross-Kerr coefficient [25].

Subgap excitation frequencies are naturally present in most superconducting resonators. Indeed, in addition to the fundamental mode, a superconducting resonator can usually support several higher harmonics with frequencies below the superconducting spectroscopic gap. For the lumped resonator design employed for the SKIDs, the fundamental resonance f_1 is of the order of 3 GHz, and, although the harmonic resonance frequencies f_n are not exactly multiples, they can be roughly estimated by $f_n = n \times f_1$ where n is an integer indexing the resonance mode. This linear relation remains valid almost all the way up to the surface plasma frequency where the dispersion relation saturates. For a low kinetic inductance superconductor (such as thin film Al, Nb, Ta, etc.) the surface plasma frequency is much larger than the superconducting spectroscopic gap. However, for a superconducting material consisting of an array of Josephson junctions, the dispersion relation saturates at the Josephson junction plasma frequency f_{JJ} , that can be smaller than the superconducting spectroscopic gap. Maleeva and coauthors model the granular aluminum (grAl) thin films as a network of Josephson junctions [25]. In the model, the Josephson

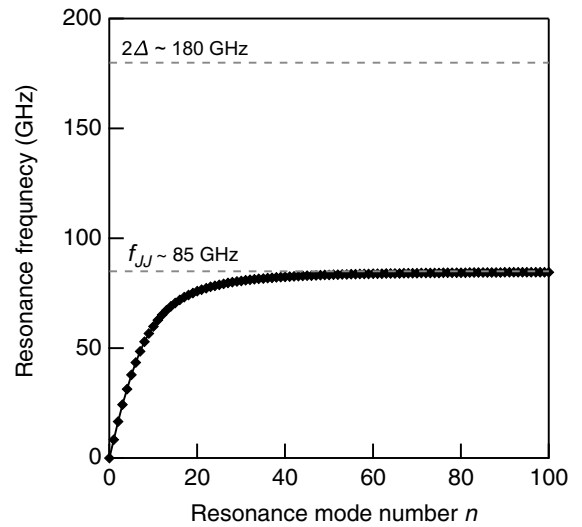


FIG. 1. Dispersion relation of the SKIDs under study. The model used to compute the dispersion relation assumes a resonator made of a superconducting material consisting of an array of Josephson junctions [25]. The dispersion relation saturates at the Josephson junction plasma frequency f_{JJ} .

junction plasma frequency is given by

$$f_{JJ} = \frac{1}{2\pi} \sqrt{\frac{2eI_c}{\hbar C_J}} \quad (2)$$

where C_J is the effective capacity between the superconducting grains. Note that the grains in the model do not necessarily correspond to the actual grains in the film, because several grains can be strongly coupled together and collectively charge as one object, as recently evidenced by scanning tunneling experiments [26]. Figure 1 shows a typical dispersion relation for a critical current density $J_c \sim 4 \text{ mA}/\mu\text{m}^2$ and $C_J \sim 40 \text{ fF}/\mu\text{m}^2$. These values are characteristics of a granular aluminum with a normal state resistivity of $\sim 900 \mu\Omega \text{ cm}$ [25]. As the density of subgap modes diverges at the plasma frequency, a strong photon absorption is expected.

In theory, the frequency selectivity is adjusted via the Josephson junction plasma frequency. In practice, a variation of the Josephson junction plasma frequency with the normal state resistivity of granular aluminum is observed [21]. Table I lists, from this previous study [21], the variation of the Josephson junction plasma frequency with the normal state resistivity. Owing to the fact that the granular aluminum film thickness is not much larger than the effective grain size, possibly giving rise to surface charging effects, and the fact that the effective intergrain capacitance likely has a nontrivial dependence on the microstructure of the film and its resistivity, the quantitative modeling of the plasma frequency as a function of thickness and resistivity is still an open question.

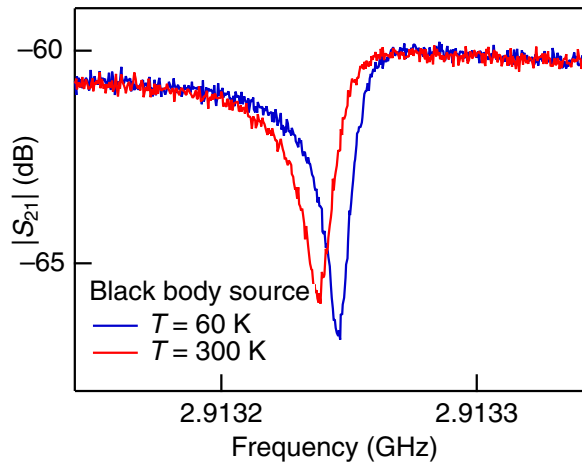


FIG. 3. Photon detection demonstration for SKID #2. The resonant frequency shifts down by 8 kHz when the temperature of a black body source illuminating the detector is increased from 60 to 300 K. The black body spectrum is filtered by a low pass with a cutoff at 110 GHz.

also be measured, as presented in the following, leading to a smaller and more realistic estimation of the change of the optical load $\Delta W_{\text{opt}}^{82-92 \text{ GHz}} \sim 0.05 \text{ pW}$. This power corresponds to imaging at 85 GHz.

Figure 4 presents measurement of the spectral response of SKID #2. The top panel shows the actual measurement: the frequency shift as a function of the optical path difference of a Martin–Puplett spectrometer [36]. The Martin–Puplett spectrometer is a Fourier transform spectrometer with a beamsplitter that consists of a grid of gold wires especially adapted for measurements in the 50 GHz–3 THz range [25,36]. The radiation source is broadband and is split into two beams and recombined with an optical path difference that is varied thanks to a moving mirror. The intensity as a function of the optical path difference, the interferogram, is the sum of interference fringes of different wavelengths. The Fourier transform of the interferogram, the spectrum, is the intensity as a function of the incoming light wavelength. For our measurements, we proceed to a lock-in detection by modulating the incoming broadband source between a black body source at nitrogen temperature (77 K) and one at room temperature. The interferogram of the SKID #2 (top panel) is almost a pure sinusoid with a few beating frequencies. The corresponding spectrum (the bottom panel) shows a maximum absorption at 86 GHz. The spectral response of the other SKIDs is similar: almost pure sinusoidal interferograms and spectra with a dominating absorption between 80 and 90 GHz. The SKIDs made out of granular aluminum with a normal state resistivity of approximately $900 \mu\Omega \text{ cm}$ are thus intrinsically selecting the 80–90 GHz frequency band, when used with a low-pass filter of 110 GHz to suppress the above gap (standard KID) optical response.

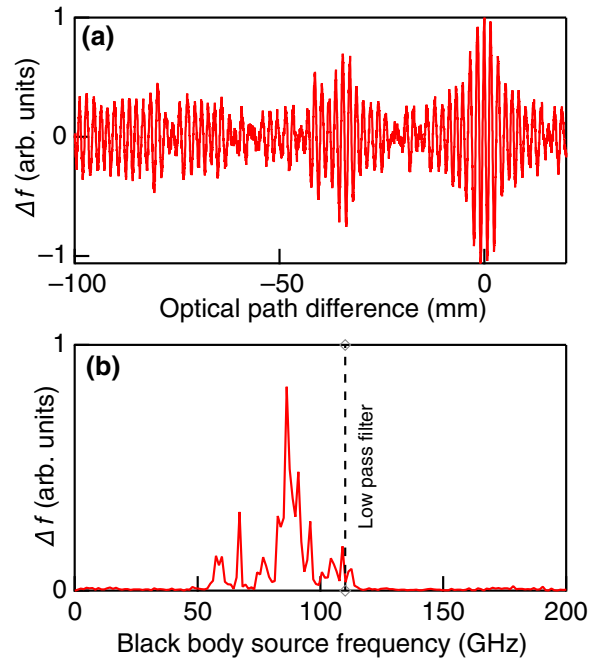


FIG. 4. Spectral response of SKID #2. (a) Frequency shift as a function of the optical path difference of a Martin–Puplett spectrometer. (b) Frequency shift as a function of incident photon frequency. Fourier transform of the top panel response.

Figure 5 shows the NSD of SKID #2. Under a constant optical load, the variation of the resonance frequency of each SKID is recorded over time and a Fourier transform is applied to obtain the NSD. At 10 Hz, the NSD of SKID #2 is $N_f \sim 4 \text{ Hz/Hz}^{0.5}$. At 10 Hz, the average NSD of the array is $28 \text{ Hz/Hz}^{0.5}$ with a standard deviation of $17 \text{ Hz/Hz}^{0.5}$. The uniformity of the array has to be improved. In particular, the response seems homogeneous, but the noise is very variable. This last point needs to be understood. Optimization of the readout power for each individual SKID may help.

Figure 6 presents the evaluation of the noise equivalent power of our detectors with Eq. (3), using the smallest value of the optical load variation. For the best detector, SKID #2, we obtain a NEP of $2.6 \times 10^{-17} \text{ W/Hz}^{0.5}$,

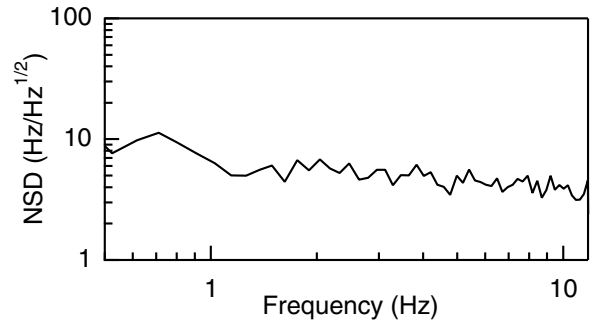


FIG. 5. Noise spectral density (NSD) of SKID #2.

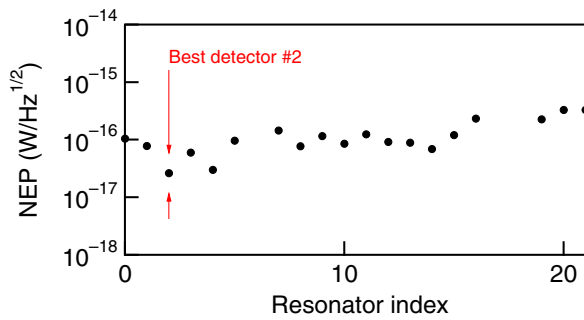


FIG. 6. Noise equivalent power of a SKIDs array for an 80–90 GHz frequency range. The SKIDs are identified by their index from 0 to 21, see Fig. 1(a) (19 out of 22 are functional resonators).

comparable to the result obtained on KIDs made of a Ti–Al bilayer operating at 100 mK for a 80–120 GHz frequency band [27]. The average NEP of the SKID array is $1.3 \times 10^{-16} \text{ W/Hz}^{0.5}$, for an operating temperature of 300 mK and a 80–90 GHz frequency band.

IV. CONCLUSION

The presented grAl SKIDs are potentially of interest for imaging applications in the 80–90 GHz band. Indeed, the upscaling of our array to a larger one is technically possible. The readout electronics are already available to address up to 400 detectors per feedline. In addition, the base temperature requirement of 300 mK, instead of 100 mK, simplifies the cryogenic setup, although the need for sub-Kelvin temperatures still implies the use of relatively complex instrumentation. The intrinsic frequency selectivity of these detectors may be used as a natural band-pass filter. However, we still have not implemented a mechanism to tune the absorption band. Further investigation may increase the required minimum operating temperature by unveiling other superconducting materials with collective subgap modes and higher critical temperatures. Moreover, THz spectroscopy on chip may be envisioned if a solution is found to tune the frequencies of the subgap modes, for instance by employing magnetic fields [37], or dc current bias.

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- [1] M. Tonouchi, Cutting-edge THz technology, *Nat. Photonics* **1**, 97 (2007).
- [2] A. Rogalski and F. Sizov, Terahertz detectors and focal plane arrays, *Opto-Electron. Rev.* **19**, 346 (2011).
- [3] S. S. Dhillon, *et al.*, The 2017 terahertz science and technology roadmap, *J. Phys. D: Appl. Phys.* **50**, 043001 (2017).
- [4] T. Encrenaz, B. Bézard, J. Crovisier, A. Coustenis, E. Lelouch, S. Gulkis, and S. Atreya, Detectability of molecular species in planetary and satellite atmospheres from their rotational transitions, *Planet. Space Sci.* **43**, 1485 (1995).
- [5] T. G. Phillips and J. Keene, Submillimeter astronomy (heterodyne spectroscopy), *Proc. IEEE* **80**, 1662 (1992).
- [6] D. Mittleman, *Sensing with Terahertz Radiation* (Springer, Berlin, 2013).
- [7] P. H. Siegel, Terahertz technology in biology and medicine, *IEEE Trans. Microw. Theory Tech.* **52**, 2438 (2004).
- [8] A. Monfardini, *et al.*, Latest NIKA results and the NIKA-2 project, *J. Low Temp. Phys.* **176**, 787 (2014).
- [9] J. Schlaerth, A. Vayonakis, P. Day, J. Glenn, J. Gao, S. Golwala, S. Kumar, H. LeDuc, B. Mazin, J. Vaillancourt, and J. Zmuidzinas, A millimeter and submillimeter kinetic inductance detector camera, *J. Low Temp. Phys.* **151**, 684 (2008).
- [10] O. Dupré, A. Benoît, M. Calvo, A. Catalano, J. Goupy, C. Hoarau, T. Klein, K. L. Calvez, B. Sacépé, A. Monfardini, and F. Levy-Bertrand, Tunable sub-gap radiation detection with superconducting resonators, *Supercond. Sci. Technol.* **30**, 045007 (2017).
- [11] P. K. Day, H. G. LeDuc, B. A. Mazin, A. Vayonakis, and J. Zmuidzinas, A broadband superconducting detector suitable for use in large arrays, *Nature* **425**, 817 (2003).
- [12] S. Doyle, P. Mauskopf, J. Naylor, A. Porch, and C. Duncombe, Lumped element kinetic inductance detectors, *J. Low Temp. Phys.* **151**, 530 (2008).
- [13] H. McCarrick, D. Flanigan, G. Jones, B. R. Johnson, P. Ade, D. Araujo, K. Bradford, R. Cantor, G. Che, P. Day, S. Doyle, H. LeDuc, M. Limon, V. Luu, P. Mauskopf, A. Miller, T. Mroczkowski, C. Tucker, and J. Zmuidzinas, Horn-coupled, commercially-fabricated aluminum lumped-element kinetic inductance detectors for millimeter wavelengths, *Rev. Sci. Instrum.* **85**, 123117 (2014).
- [14] A. Patel, A.-D. Brown, W.-T. Hsieh, T. Stevenson, S. H. Moseley, K. U-yen, N. Ehsan, E. Barrentine, G. Manos, and E. J. Wollack, Fabrication of MKIDS for the MicroSpec spectrometer, *IEEE Trans. Appl. Supercond.* **23**, 2400404 (2013).
- [15] T. Matsumura, *et al.*, Mission design of LiteBIRD, *J. Low Temp. Phys.* **176**, 733 (2014).
- [16] M. Griffin, *et al.*, in *Millimeter, Submillimeter, and Far-Infrared Detectors and Instrumentation for Astronomy VIII* (SPIE Astronomical Telescopes + Instrumentation, Edinburgh, United Kingdom, 2016), Vol. 9914, p. 991407.
- [17] J. Baselmans, J. Bueno, O. Yurduseven, S. Yates, N. Llombart, V. Murugesan, D. J. Thoen, A. Baryshev, and A.

- Neto, in *41st Int. Conf. on Infrared, Millimeter, and Terahertz waves (IRMMW-THz)* (IEEE, Copenhagen, Denmark, 2016).
- [18] S. Rowe, E. Pascale, S. Doyle, C. Dunscombe, P. Hargrave, A. Papageorgio, K. Wood, P. A. R. Ade, P. Barry, A. Bideaud, T. Brien, C. Dodd, W. Grainger, J. House, P. Mauskopf, P. Moseley, L. Spencer, R. Sudiwala, C. Tucker, and I. Walker, A passive terahertz video camera based on lumped element kinetic inductance detectors, *Rev. Sci. Instrum.* **87**, 033105 (2016).
- [19] <https://sequestim.com>
- [20] R. Adam, *et al.*, The NIKA2 large-field-of-view millimetre continuum camera for the 30 m IRAM telescope, *A&A* **609**, A115 (2018).
- [21] F. Levy-Bertrand, T. Klein, T. Grenet, O. Dupré, A. Benoît, A. Bideaud, O. Bourrion, M. Calvo, A. Catalano, A. Gomez, J. Goupy, L. Grünhaupt, U. v. Luepke, N. Maleeva, F. Valenti, I. M. Pop, and A. Monfardini, Electrodynamics of granular aluminum from superconductor to insulator: Observation of collective superconducting modes, *Phys. Rev. B* **99**, 094506 (2019).
- [22] P. G. D. Gennes, *Superconductivity of Metals And Alloys*. Chap. 6.5 (CRC Press, Boca Raton, 2018).
- [23] L. J. Swenson, P. K. Day, B. H. Eom, H. G. Leduc, N. Llobart, C. M. McKenney, O. Noroozian, and J. Zmuidzinas, Operation of a titanium nitride superconducting microresonator detector in the nonlinear regime, *J. Appl. Phys.* **113**, 104501 (2013).
- [24] A. Kher, P. K. Day, B. H. Eom, J. Zmuidzinas, and H. G. Leduc, Kinetic inductance parametric up-converter, *J. Low Temp. Phys.* **184**, 480 (2016).
- [25] N. Maleeva, L. Grünhaupt, T. Klein, F. Levy-Bertrand, O. Dupre, M. Calvo, F. Valenti, P. Winkel, F. Friedrich, W. Wernsdorfer, A. V. Ustinov, H. Rotzinger, A. Monfardini, M. V. Fistul, and I. M. Pop, Circuit quantum electrodynamics of granular aluminum resonators, *Nat. Commun.* **9**, 3889 (2018).
- [26] F. Yang, T. Gozliniski, T. Storbeck, L. Grünhaupt, I. M. Pop, and W. Wulfhchel, Microscopic charging and in-gap states in superconducting granular aluminum, *Phys. Rev. B* **102**, 104502 (2020).
- [27] A. Catalano, J. Goupy, H. le Sueur, A. Benoit, O. Bourrion, M. Calvo, A. D'addabbo, L. Dumoulin, F. Levy-Bertrand, J. Macías-Pérez, S. Marnieros, N. Ponthieu, and A. Monfardini, Bi-layer kinetic inductance detectors for space observations between 80–120 GHz, *Astron. Astrophys.* **580**, A15 (2015).
- [28] R. W. Cohen and B. Abeles, Superconductivity in granular aluminum films, *Phys. Rev.* **168**, 444 (1968).
- [29] R. H. Parmenter, Isospin formulation of the theory of a granular superconductor, *Phys. Rev.* **154**, 353 (1967).
- [30] G. Deutscher, H. Fenichel, M. Gershenson, E. Grünbaum, and Z. Ovadyahu, Transition to zero dimensionality in granular aluminum superconducting films, *J. Low Temp. Phys.* **10**, 231 (1973).
- [31] R. C. Dynes and J. P. Garno, Metal-Insulator Transition in Granular Aluminum, *Phys. Rev. Lett.* **46**, 137 (1981).
- [32] U. S. Pracht, N. Bachar, L. Benfatto, G. Deutscher, E. Farber, M. Dressel, and M. Scheffler, Enhanced Cooper pairing versus suppressed phase coherence shaping the superconducting dome in coupled aluminum nanograins, *Phys. Rev. B* **93**, 100503(R) (2016).
- [33] The optical filter stack consists of two thermal filters at the 150 K and the 50 K stages and five low passes filters: a 18 cm^{-1} filter at the 50 K stage, a 13.5 cm^{-1} one at 4 K, a 12 cm^{-1} one at 1 K, and two filters of 5.7 and 3.9 cm^{-1} at 300 mK.
- [34] A. Catalano, A. Bideaud, O. Bourrion, M. Calvo, A. Fasano, J. Goupy, F. Levy-Bertrand, J. F. Macias-Perez, N. Ponthieu, Q. Y. Tang, and A. Monfardini, Sensitivity of LEKID for space applications between 80 and 600 GHz, *A&A* **641**, A179 (2020).
- [35] F. Valenti, F. Henriques, G. Catelani, N. Maleeva, L. Grünhaupt, U. von Lüpke, S. T. Skacel, P. Winkel, A. Bilmes, A. V. Ustinov, J. Goupy, M. Calvo, A. Benoît, F. Levy-Bertrand, A. Monfardini, and I. M. Pop, Interplay between Kinetic Inductance, Nonlinearity, and Quasiparticle Dynamics in Granular Aluminum Microwave Kinetic Inductance Detectors, *Phys. Rev. Appl.* **11**, 054087 (2019).
- [36] D. Martin and E. Pulett, Polarised interferometric spectrometry for the millimetre and submillimetre spectrum, *Infrared Phys.* **10**, 105 (1970).
- [37] K. Borisov, D. Rieger, P. Winkel, F. Henriques, F. Valenti, A. Ionita, M. Wessbecher, M. Spiecker, D. Gusenkova, I. M. Pop, and W. Wernsdorfer, Superconducting granular aluminum resonators resilient to magnetic fields up to 1 Tesla, *Appl. Phys. Lett.* **117**, 120502 (2020).